

Electromagnetic Theory I

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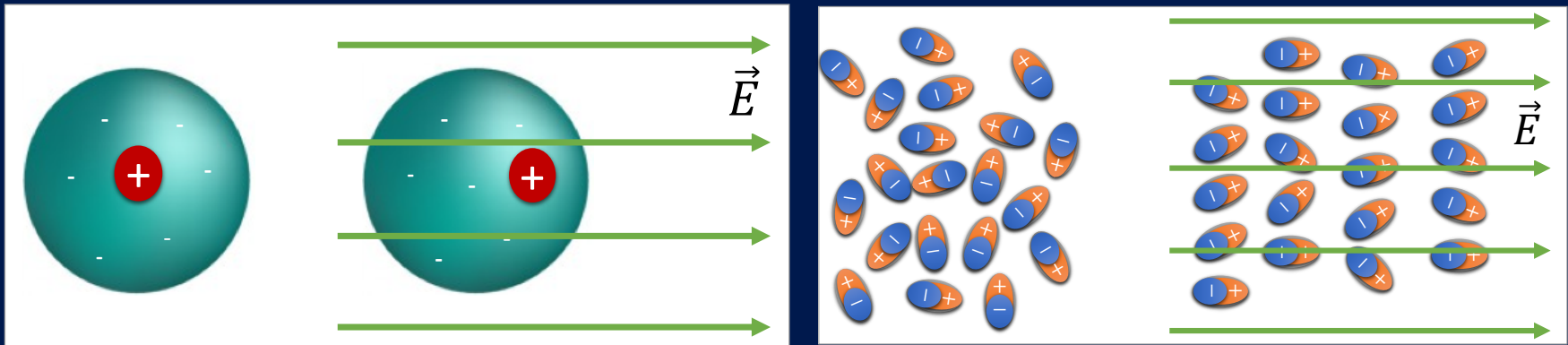
Chapter 4: Electric Fields in Matter

- * Polarization
- * The Field of a Polarized Object
- * The Electric Displacement
- * Linear Dielectrics

4.1 Polarization

Dielectrics (e.g. plastic, glass, ...)

- In the most materials, charges are linked to atoms or molecules.
- They are considered bound and only can move within the atom or molecule.
- These restricted microscopic displacements give rise for the physical properties of dielectric materials.
- Electric fields can affect a charge distribution of a dielectric atom or molecule by *stretching and rotating* them.



4.1 Polarization

Induced Dipoles

What happens to a neutral atom when it is placed in an electric field?

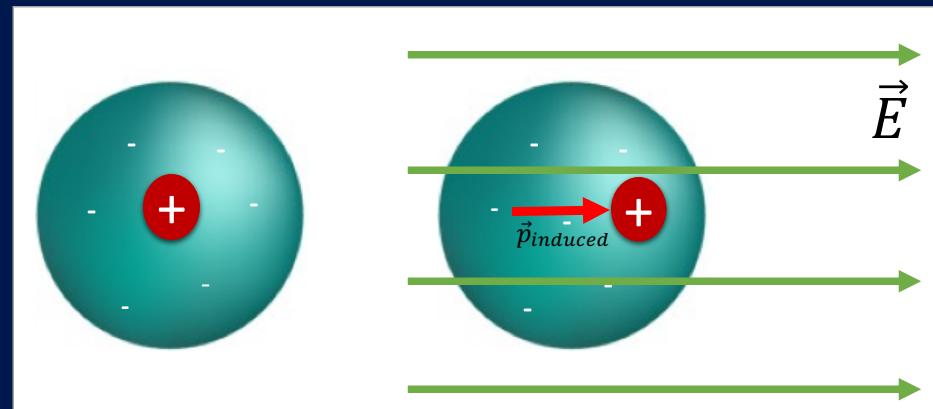
- The nucleus is pushed in the direction of the field, and the electrons the opposite way.
- The two opposing forces on nucleolus reach a balance, leaving the atom polarized.
- The atom acquires a tiny dipole moment \vec{p} , which points in the same direction as \vec{E} .
- This induced dipole moment is approximately proportional to the field (as long as the latter is not too strong)

$$\vec{p} = \alpha \vec{E}$$

α is called **atomic polarizability**

Atomic Polarizabilities ($\alpha/4\pi\epsilon_0$) in units of $10^{-30}m$

H	He	Li	Be	C	Ne	Na	Ar	K	Cs
0.667	0.205	24.3	5.60	1.76	0.396	24.1	1.64	43.4	59.6



4.1 Polarization

Example: An atom consists of a point nucleus ($+Q$) surrounded by a cloud ($-Q$) of radius a . Calculate the atomic polarizability of such an atom.
uniformly charged spherical

■ As simple model, we will assume

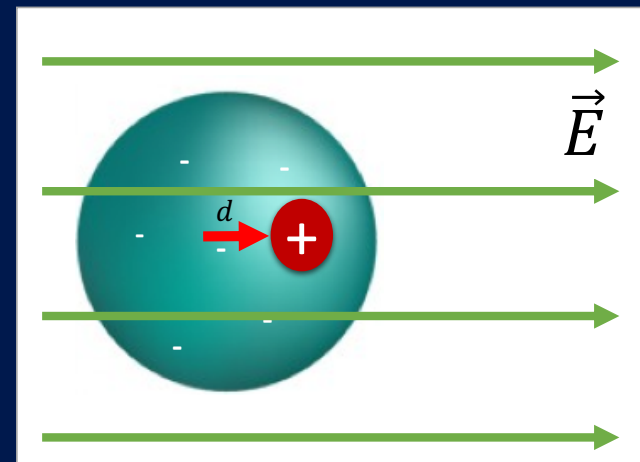
- The electron cloud retains its spherical shape.
- Equilibrium occurs when the nucleus is displaced a distance d from the center of the sphere and pulling by the internal field produced by the electron cloud.

At equilibrium,

$$E = E_e = \frac{1}{4\pi\epsilon_0} \frac{Qd}{a^3} = \frac{1}{4\pi\epsilon_0} \frac{p}{a^3}$$

$$\rightarrow p = 4\pi\epsilon_0 a^3 E$$

$$\rightarrow \alpha = 4\pi\epsilon_0 a^3 \rightarrow \alpha = 3\epsilon_0 v$$



4.1 Polarization

- **Example:** According to quantum mechanics, the electron cloud for a hydrogen atom in the ground state has a charge density:

$$\rho(\vec{r}) = \frac{q}{\pi a^3} e^{-2\frac{r}{a}}$$

where q is the charge of the electron and a is the Bohr radius. Find the atomic polarizability of such an atom.

- Using Gauss's Law, the electric field due to the electron cloud is given by

$$E_e = \frac{1}{4\pi\epsilon_0} \frac{Q_{enc}}{r^2}$$

$$Q_{enc} = \int_0^r \frac{q}{\pi a^3} e^{-2\frac{r'}{a}} 4\pi r'^2 dr' = q \left[1 - e^{-2\frac{r}{a}} \left(1 + 2\frac{r}{a} + 2\frac{r^2}{a^2} \right) \right]$$

4.1 Polarization

- Similar to the previous example, the proton will be shifted a distance d at equilibrium, and hence

$$E = E_e = \frac{1}{4\pi\epsilon_0} \frac{q \left[1 - e^{-2\frac{d}{a}} \left(1 + 2\frac{d}{a} + 2\frac{d^2}{a^2} \right) \right]}{d^2}$$

For $d \ll a$

$$\begin{aligned} \left[1 - e^{-2\frac{d}{a}} \left(1 + 2\frac{d}{a} + 2\frac{d^2}{a^2} \right) \right] &= 1 - \left(1 - 2\frac{d}{a} + 2\frac{d^2}{a^2} - \frac{4}{3}\frac{d^3}{a^3} + \dots \right) \left(1 + 2\frac{d}{a} + 2\frac{d^2}{a^2} \right) \\ &= \frac{4}{3}\frac{d^3}{a^3} + \text{HOT} \end{aligned}$$

$$E = \frac{1}{3\pi\epsilon_0} \frac{qd}{a^3} = \frac{1}{3\pi\epsilon_0} \frac{p}{a^3}$$

$$\rightarrow p = 3\pi\epsilon_0 a^3 E$$

$$\rightarrow \alpha = 3\pi\epsilon_0 a^3$$

4.1 Polarization

Polarizability Tensor

- For example in Carbon dioxide, when the field is at some angle to the axis, it will have parallel and perpendicular components,

$$\vec{p} = \alpha_{\parallel} \vec{E}_{\parallel} + \alpha_{\perp} \vec{E}_{\perp}$$

in general, linear response theory gives

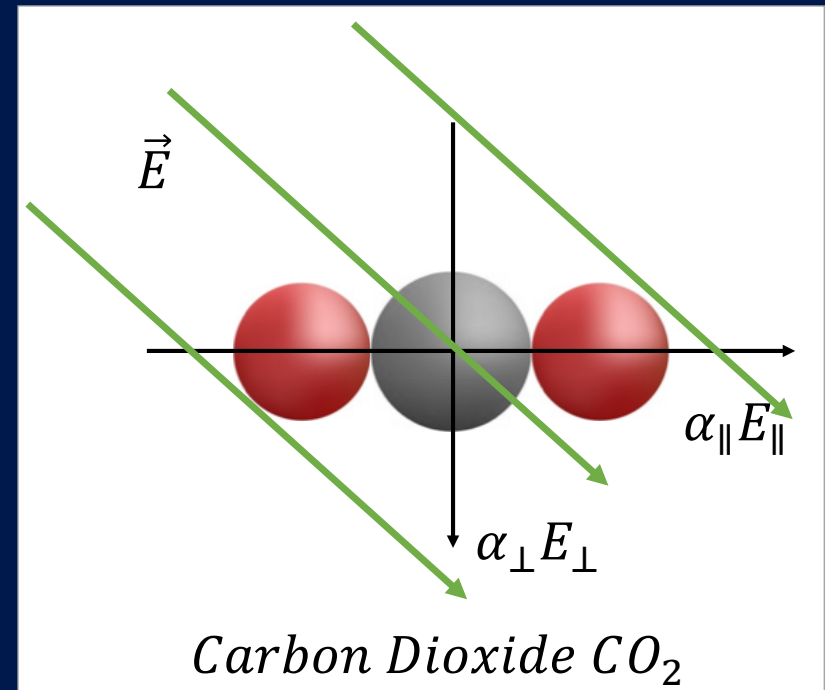
$$\vec{p} = \bar{\alpha} \vec{E}$$

$$p_x = \alpha_{xx} E_x + \alpha_{xy} E_y + \alpha_{xz} E_z$$

$$p_y = \alpha_{yx} E_x + \alpha_{yy} E_y + \alpha_{yz} E_z$$

$$p_z = \alpha_{zx} E_x + \alpha_{zy} E_y + \alpha_{zz} E_z$$

$\bar{\alpha}$ is called **polarizability tensor**

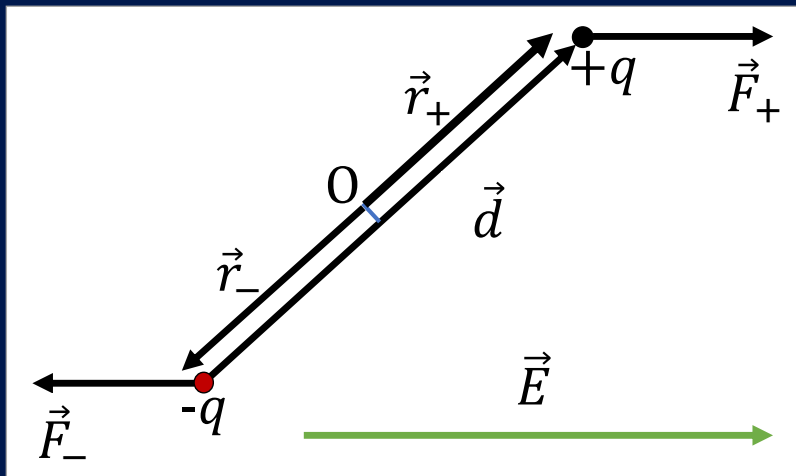


4.1 Polarization

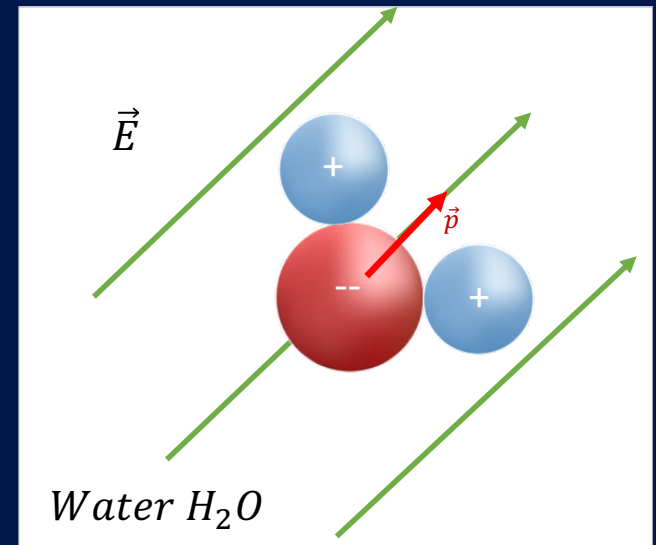
Alignment of Polar Molecules

- Polar molecules such as water have intrinsic dipole moment, hence when an external field is applied, there will be a net torque on the molecule. It will be free to rotate and swing around until it points in the direction of the applied field.

If the electric field is uniform



$$\vec{F}_{net} = 0$$



$$\vec{N} = \vec{r}_+ \times \vec{F}_+ + \vec{r}_- \times \vec{F}_- = \frac{\vec{d}}{2} \times q\vec{E} + \frac{-\vec{d}}{2} \times -q\vec{E} = q\vec{d} \times \vec{E} = \vec{p} \times \vec{E}$$

4.1 Polarization

If the electric field is not uniform

$$\begin{aligned}\vec{F}_{net} &= \vec{F}_+ + \vec{F}_- = q(\vec{E}_+ - \vec{E}_-) = q\Delta\vec{E} \approx q\vec{\nabla}\vec{E} \cdot \vec{d} \\ &= q(\vec{d} \cdot \vec{\nabla})\vec{E} = (\vec{p} \cdot \vec{\nabla})\vec{E}\end{aligned}$$

$$\vec{N} = \vec{p} \times \vec{E} + \vec{r} \times \vec{F}_{net} = \vec{p} \times \vec{E} + \vec{r} \times (\vec{p} \cdot \vec{\nabla})\vec{E}$$

4.1 Polarization

Energy of an Ideal Dipole

- For a physical dipole,

$$U = \sum_{n=1}^2 q_i V(\vec{r}_i) = qV(\vec{r}_+) - qV(\vec{r}_-) = qV\left(\vec{r} + \frac{\vec{d}}{2}\right) - qV\left(\vec{r} - \frac{\vec{d}}{2}\right) = q \left(- \int_{\vec{r} - \frac{\vec{d}}{2}}^{\vec{r} + \frac{\vec{d}}{2}} \vec{E} \cdot d\vec{l} \right)$$

- For an ideal dipole, $\vec{d} \rightarrow 0$

$$U = \lim_{\vec{d} \rightarrow 0} q \left(- \int_{\vec{r} - \frac{\vec{d}}{2}}^{\vec{r} + \frac{\vec{d}}{2}} \vec{E} \cdot d\vec{l} \right) = -q\vec{E} \cdot \left[\left(\vec{r} + \frac{\vec{d}}{2} \right) - \left(\vec{r} - \frac{\vec{d}}{2} \right) \right] = -q\vec{E} \cdot \vec{d} = -\vec{p} \cdot \vec{E}$$

4.1 Polarization

Example: A dipole \vec{p} is a distance \vec{r} from a point charge q , and oriented so that \vec{p} makes an angle θ_0 with the vector \vec{r} from q to \vec{p} . What is the force on \vec{p} ?

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \frac{q}{r^2} \hat{r} = \frac{1}{4\pi\epsilon_0} \frac{q}{(x^2 + y^2 + z^2)^{\frac{3}{2}}} (x\hat{x} + y\hat{y} + z\hat{z}) =$$

$$\vec{F} = (\vec{p} \cdot \nabla) \vec{E}$$

$$\begin{aligned} F_x &= \left(p_x \frac{\partial}{\partial x} + p_y \frac{\partial}{\partial y} + p_z \frac{\partial}{\partial z} \right) E_x \\ &= \frac{q}{4\pi\epsilon_0} \left[p_x \frac{-2x^2 + y^2 + z^2}{(x^2 + y^2 + z^2)^{\frac{5}{2}}} - p_y \frac{-3xy}{(x^2 + y^2 + z^2)^{\frac{3}{2}}} - p_z \frac{-3xz}{(x^2 + y^2 + z^2)^{\frac{3}{2}}} \right] = \frac{q}{4\pi\epsilon_0} \left[\frac{\vec{p}}{r^3} - \frac{3\vec{r}(\vec{p} \cdot \vec{r})}{r^5} \right]_x \\ \vec{F} &= \frac{q}{4\pi\epsilon_0} \left[\frac{\vec{p}}{r^3} - \frac{3\vec{r}(\vec{p} \cdot \vec{r})}{r^5} \right] = \frac{q}{4\pi\epsilon_0 r^3} [\vec{p} - 3\hat{r}(\vec{p} \cdot \hat{r})] = \frac{q}{4\pi\epsilon_0 r^3} [\vec{p} - 3\hat{r} p \cos \theta_0] \end{aligned}$$

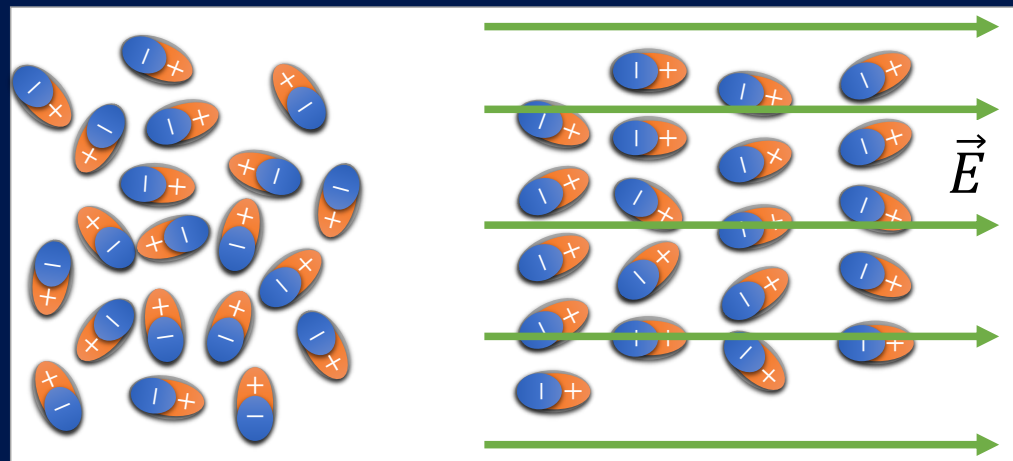
4.1 Polarization

When a piece of dielectric material is placed in an electric field, its constituent dipoles tend to point along the direction of the field and the material become **polarized**

This Polarization will change the electric strength in a dielectric medium which can be quantified through the polarization vector (\vec{P})

\vec{P} is the total dipole moment per unit volume

$$\vec{P} = \lim_{v \rightarrow 0} \frac{1}{v} \sum_{\vec{p}_i \in v} \vec{p}_i$$



4.2 The Field of a Polarized Object

Bound Charges

For a single dipole,
$$V_{dip}(\vec{r}) = \frac{1}{4\pi\epsilon_0} \frac{\vec{p} \cdot \hat{r}}{r^2}$$

For a collection of dipoles represented in the continuous limit by the polarization vector, the potential generated by the microscopic dipoles is given by

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\vec{P}(\vec{r}') \cdot \hat{r}}{r^2} d\tau' \rightarrow V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \vec{P}(\vec{r}') \cdot \vec{\nabla}' \left(\frac{1}{r} \right) d\tau'$$

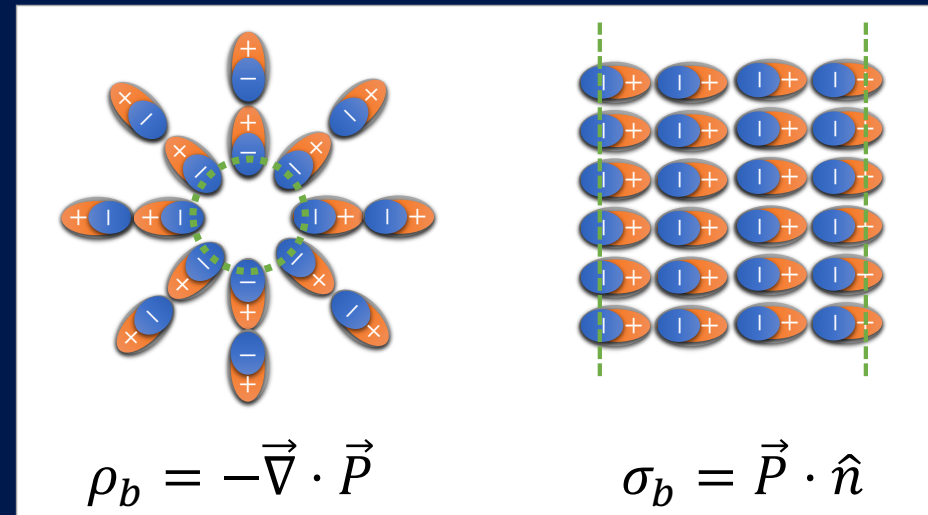
$$\vec{\nabla}' \left(\frac{1}{r} \right) = \frac{\hat{r}}{r^2}$$

4.2 The Field of a Polarized Object

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \int \vec{\nabla}' \cdot \left(\frac{\vec{P}(\vec{r}')}{r} \right) - \left(\frac{\vec{\nabla}' \cdot \vec{P}(\vec{r}')}{r} \right) d\tau'$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \oint \frac{\vec{P}(\vec{r}')}{r} \cdot d\vec{a}' + \frac{1}{4\pi\epsilon_0} \int - \left(\frac{\vec{\nabla}' \cdot \vec{P}(\vec{r}')}{r} \right) d\tau'$$

$$V(\vec{r}) = \frac{1}{4\pi\epsilon_0} \oint \frac{\sigma_b}{r} \cdot da' + \frac{1}{4\pi\epsilon_0} \int \frac{\rho_b}{r} d\tau'$$



4.2 The Field of a Polarized Object

Example: Find the electric field produced by a uniformly polarized sphere of radius R .

$$\begin{aligned}
 P(\vec{r}) &= P_0 \hat{z} \\
 \rho_b &= -\vec{\nabla} \cdot \vec{P} = 0 \\
 \sigma_b &= \vec{P} \cdot \hat{r} = P_0 \cos \theta
 \end{aligned}
 \rightarrow
 \begin{aligned}
 V(r, \theta) &= \begin{cases} \frac{P_0 r}{3\epsilon_0} \cos \theta & r \leq R \\ \frac{P_0 R^3}{3\epsilon_0 r^2} \cos \theta & r \geq R \end{cases}
 \end{aligned}$$

Outside the sphere the potential is identical to that of a perfect dipole at the origin,

$$\vec{E}(r, \theta) = -\vec{\nabla}V = \begin{cases} -\frac{\vec{p}}{4\pi\epsilon_0 R^3} & r \leq R \\ \frac{1}{4\pi\epsilon_0 r^3} (3\vec{p} \cdot \hat{r} \hat{r} - \vec{p}) & r \geq R \end{cases}$$

$$\vec{p} = \frac{4}{3} \pi R^3 \vec{P}$$

