Electromagnetic Theory I

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May 19, 2021

Chapter 5: Magnetostatics

*** The Lorentz Force Law *** The Biot-Savart Law** \cdot The Divergence and Curl of \overrightarrow{B} *** Magnetic Vector Potential (** \vec{A} **)**

Introduction

Lodestone (Magnetite $Fe₃O₄$) Known by Ancient civilizations

³ **Chapter 5: Magnetostatics 5.1 The Lorentz Force Law**

n Orsted 1820 **Battery** Win Wooden
Wire Clamp Compass Current Flow Current Flow **Battery Battery** Wire Wire Wooden Wooden **Wire Clamp Wire Clamp Compass Compass**

Magnetic Fields due to currents

Moving charges generate magnetic fields

Force between two wires

The forces between charges in motion:

- Attraction of parallel currents
- Repulsion of antiparallel currents
- Does charging up the wires make simply the electrical repulsion
- of like charges?
- The wires are electrically neutral.
- It is not electrostatic in nature.

Lorentz Force Law

The net force on Q in the presence of both electric and magnetic fields:

$$
\vec{F} = Q\big(\vec{E} + \vec{v}{\times}\vec{B}\big)
$$

The magnetic force in a charge $Q,$ moving with velocity $\vec{\nu}$ in a magnetic field \vec{B}_{\parallel}

$$
\vec{F}_{mag} = Q\big(\vec{\nu}{\times}\vec{B}\big)
$$

Magnetic forces do not work!

$$
dW_{mag} = \vec{F}_{mag} \cdot d\vec{l} = Q(\vec{v} \times \vec{B}) \cdot \vec{v} dt = 0
$$

Magnetic forces may change the direction in which a particle moves, but they cannot speed it up or slow it down.

Example: Cyclotron motion?: If a charge Q moves counterclockwise, with speed ν , around a circle of radius R, in a plane perpendicular to \vec{B} . what path will it follow?

Cyclotrons represent the first model of modern particle accelerators. The magnetic force is inward, and has a fixed magnitude to sustain uniform circular motion:

$$
|\vec{F}_{mag}| = |Q(\vec{v} \times \vec{B})| = QvB = \frac{mv^2}{R} \to \omega = \frac{v}{R} = \frac{QB}{m}
$$

 $\overline{p} = mv = QBR$ Experimentally the momentum of a particle can be found using

If the velocity has a parallel component to the magnetic field, we get helical motion

Example: Cycloid motion?: If a charge Q at rest is released from the origin, what path it will follow?

 $\vec{B} = B\hat{x}$ $\vec{v} = \dot{y}\hat{y} + \dot{z}\hat{z}$ $\vec{F} = Q(\vec{E} + \vec{v} \times \vec{B})$ $\vec{E} = E\hat{z}$ $\vec{F} = QE\hat{z} + Q$ $\widehat{\chi}$ \widehat{y} \widehat{z} $0 \quad \dot{y} \quad \dot{z}$ " 0 0 $\vec{F} = Q(B\dot{z}\hat{y} + \overline{(E-B\dot{y})}\hat{z})$ $\vec{v}_0=0$, $\vec{r}_0=0$

$$
\vec{F} = Q(B\dot{z}\hat{y} + (E - B\dot{y})\hat{z}) = m\vec{a} = m(\dot{y}\hat{y} + \ddot{z}\hat{z})
$$

$$
Q(E - B\dot{y}) = m\ddot{z}
$$

$$
\ddot{y} = \frac{QB}{m}\ddot{z} = \frac{Q^2B^2}{m^2}\left(\frac{E}{B} - \dot{y}\right) = \omega^2\left(\frac{E}{B} - \dot{y}\right)
$$

$$
QB\dot{z} = m\ddot{y}
$$

$$
\ddot{z} = -\frac{QB}{m}\ddot{y} = -\frac{Q^2B^2}{m^2}\dot{z} = -\omega^2\dot{z}
$$

$$
z(t) = -C_1\sin\omega t + C_2\cos\omega t + C_4
$$

$$
y(t) = -\frac{E}{\omega B} \sin \omega t + \frac{E}{B} t
$$

$$
z(t) = -\frac{E}{\omega B} \cos \omega t + \frac{E}{\omega B}
$$

 $\vec{v}_0 = 0$, $\vec{r}_0 = 0 \rightarrow$

$$
let R = \frac{E}{\omega B}
$$

 $(z - R)^2 = (-R \cos \omega t + R - R)^2 = R^2 \cos^2 \omega t$ $(y - R\omega t)^2 = (R \sin \omega t + R\omega t - R \omega t)^2 = R^2 \sin^2 \omega t$

$$
(y - R\omega t)^2 + (z - R)^2 = R^2 \sin^2 \omega t + R^2 \cos^2 \omega t = R^2
$$

Cycloid motion: a circle, of radius R , with its center $(0, R\omega t, R)$ travels in the y-direction at a constant speed, $v = E/B$

Currents

A line charge λ traveling down at velocity $\vec{\mathcal{v}}$

 $\vec{l} = \lambda \vec{v}$

The magnetic force on a segment of current-carrying wire is

$$
\vec{F}_{mag} = \int \vec{v} \times \vec{B} \, dq = \int \vec{v} \times \vec{B} \, \lambda dl = \int \vec{I} \times \vec{B} \, dl
$$

 \vec{l} and $d\vec{l}$ are in the same direction

For a steady current along the wire

$$
\vec{F}_{mag} = I \int d\vec{l} \times \vec{B}
$$

When charge flows over surface with surface current density $\vec{\kappa}$

$$
\vec{\kappa} = \frac{d\vec{l}}{dl^{\perp}} = \sigma \vec{\nu}
$$

The magnetic force on the sheet $\vec{\kappa}$:current per unit width-perpendicular-to-flow

$$
\vec{F}_{mag} = \int \vec{v} \times \vec{B} \, dq = \int \vec{v} \times \vec{B} \, \sigma da = \int \vec{\kappa} \times \vec{B} da
$$

When charge flows through a volume with volume current density \vec{J}

$$
\vec{J} = \frac{d\vec{l}}{da^{\perp}} = \rho \vec{\nu}
$$

The magnetic force on the sheet $\vec{J}:$ $\vec{J$

$$
\vec{F}_{mag} = \int \vec{v} \times \vec{B} \, dq = \int \vec{v} \times \vec{B} \, \rho d\tau = \int \vec{J} \times \vec{B} \, d\tau
$$

Continuity equation

 \dot{J} $\vec{J} =$ $d\vec{l}$ da^{\perp}

The current crossing a surface S can be written as

be written as
\n
$$
I = \int_{S} Jda_{\perp} = \int_{S} \vec{J} \cdot d\vec{a}
$$

But, the total charge per unit time leaving a volume V is

Conservation of
charge

$$
\int_{S} \vec{f} \cdot d\vec{a} = \int_{V} \vec{\nabla} \cdot \vec{f} d\tau = -\frac{d}{dt} \int_{V} \rho d\tau = -\int_{V} \frac{\partial \rho}{\partial t} d\tau
$$

Continuity equation

$$
\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0
$$

(Local charge conservation)

$$
\frac{\partial \rho}{\partial t} = 0 \to \vec{\nabla} \cdot \vec{j} = 0 \to \int_{S} \vec{j} \cdot d\vec{a} = 0 \to \sum_{k} I_k = 0 \qquad \text{(Kirchhoff's current law)}
$$

Example: Suppose the current density distributed over a wire of circular cross section, with radius α is given by $J = ks$. Find the total current in the wire

$$
I = \int_{S} \vec{J} \cdot d\vec{a} = \int_{0}^{2\pi} \int_{0}^{a} ks \, s ds \, d\phi = 2\pi k \int_{0}^{a} s^2 ds = \frac{2\pi k a^3}{3}
$$

